



The Abdus Salam  
**International Centre  
for Theoretical Physics**

Quantum Fluids, Quantum Field  
Theory and Gravity,  
Chernogolovka  
18 October 2019

# LOCALIZATION AND ERGODIC TRANSITIONS IN LOG-NORMAL ROSENZWEIG-PORTER RANDOM MATRIX ENSEMBLE.

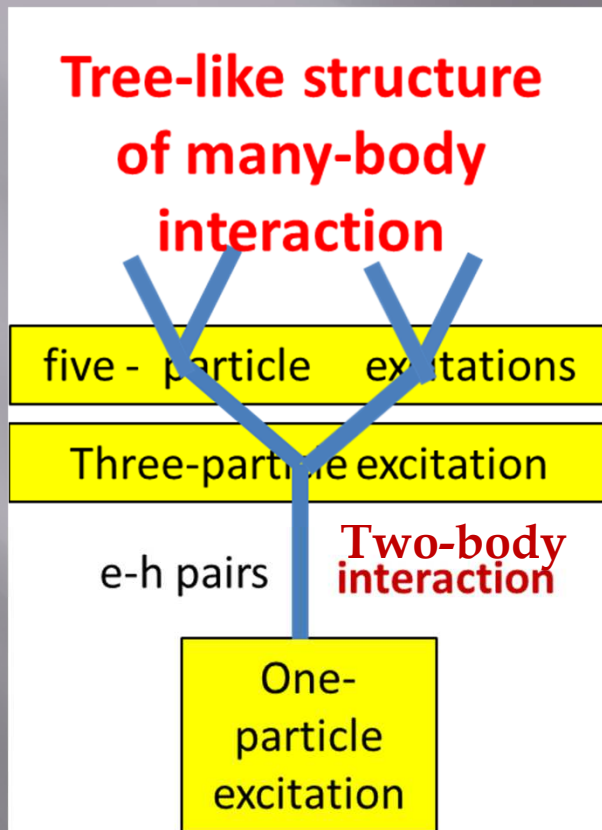
V.E.Kravtsov  
ICTP, Trieste

## Collaboration:

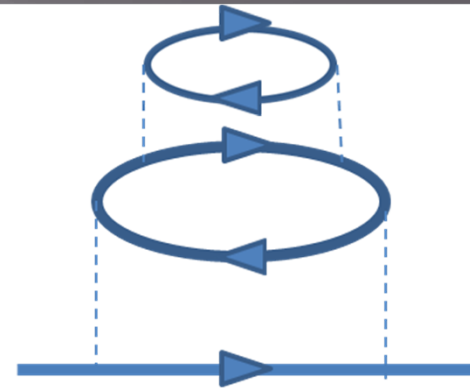
Boris Altshuler (Columbia Uni.)  
Emilio Cuevas (Murcia)  
Lev Ioffe (Google)  
Ivan Khaymovich, (MPI, Dresden)



# MBL and sing-particle localization on hierarchical graphs



Altshuler, Gefen, Kamenev, Levitov, 1997



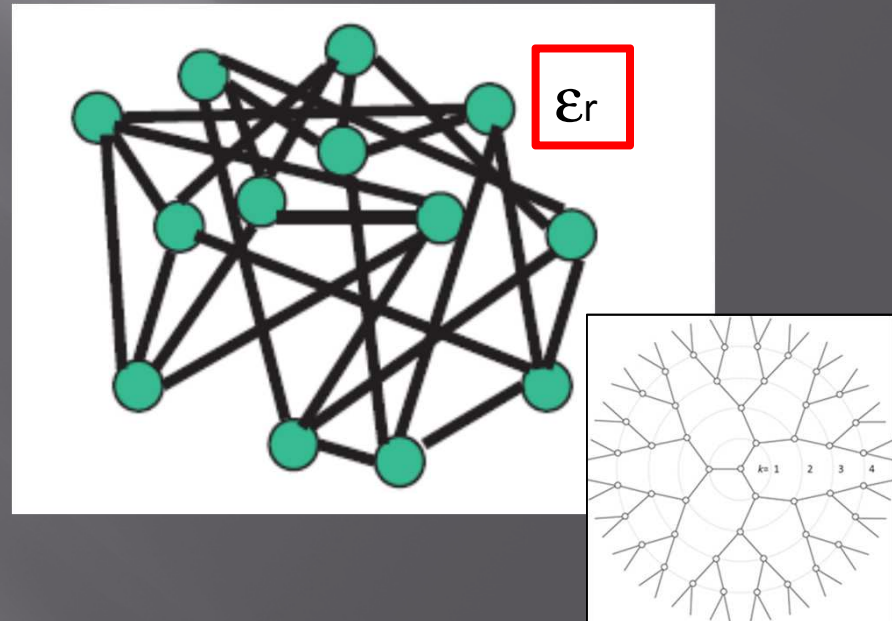
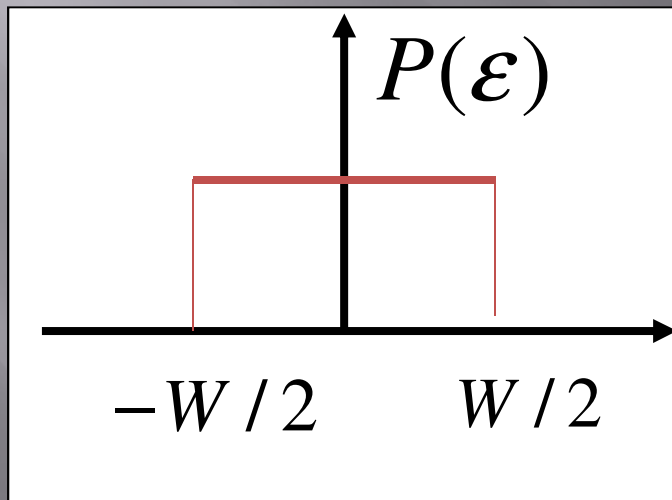
Basko, Aleiner, Altshuler, 2005



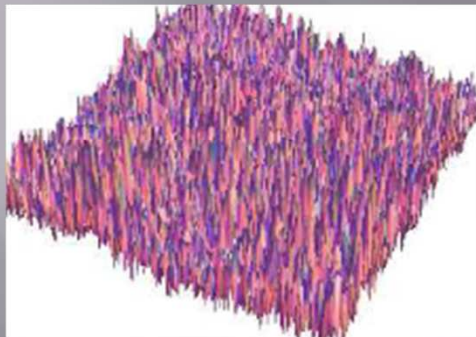
# Anderson model on hierarchical boundary-less graphs

$$H = -I \sum_{\langle r, r' \rangle} c_r^+ c_{r'} + \sum_r \varepsilon_r n_r$$

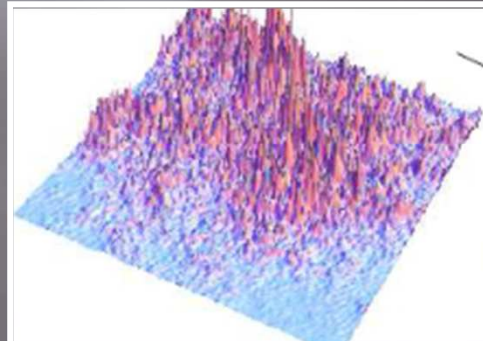
Disorder strength  $W$



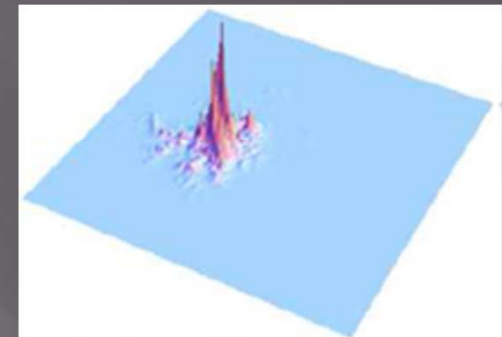
# Localization and ergodic transitions



Ergodic extended (EE)



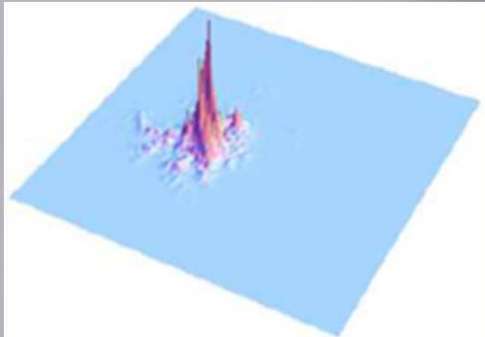
Extended but not ergodic (NEE)



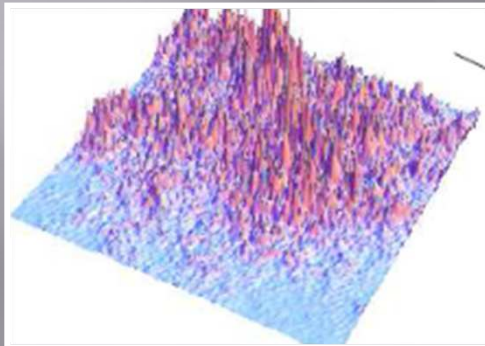
Localized (LOC)

in a  
**3D Anderson model:  
NEE states only at  
the AT point.**  
ters.

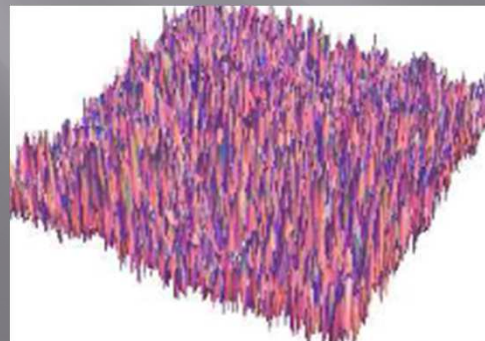
# LOCALIZED, EXTENDED ERGODIC AND EXTENDED NON-ERGODIC PHASES



Finite number of occupied sites  
in thermodynamic limit



Infinite number of occupied sites  
but zero fraction of all sites  
in the thermodynamic limit



Finite fraction of occupied sites

# Multifractal NEE

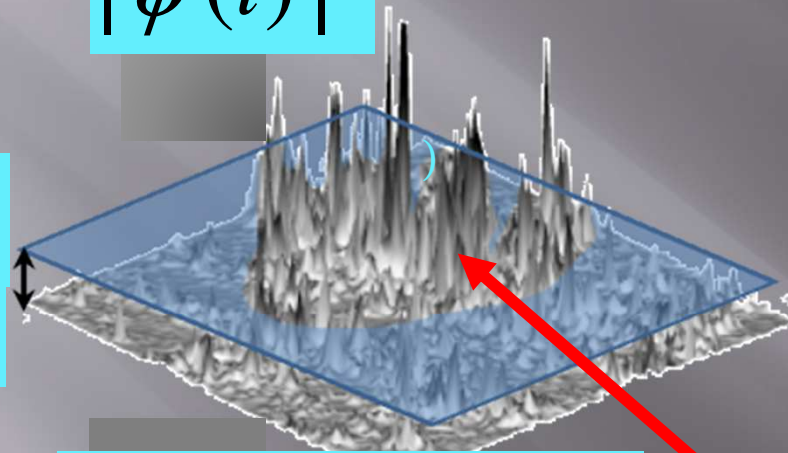
$$|\psi(i)|^2$$

Map of the regions with amplitude larger than the chosen level

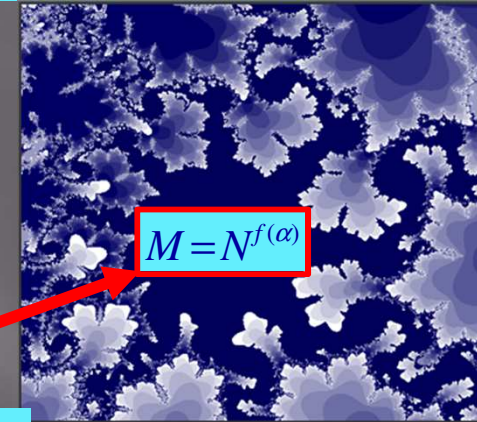
$$N^{-\alpha}$$

$$|\psi_i|^2 \leq 1$$

$$\alpha \geq 0$$



$N$  sites  $\{i\}$  in the sample



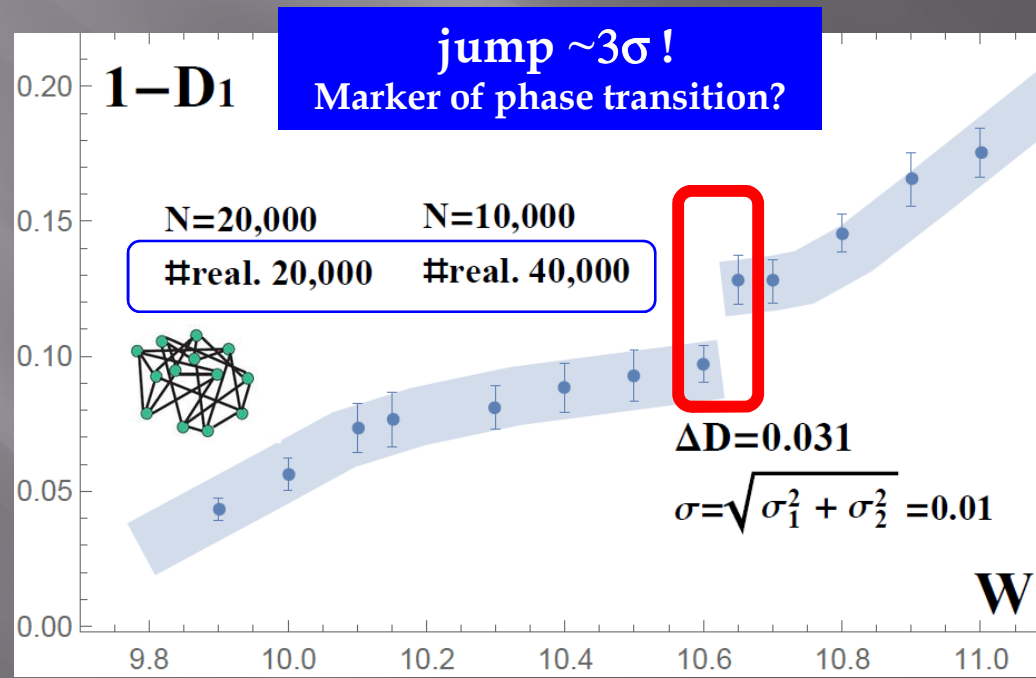
$M$  sites in this region

$$\sum_i |\psi(i)|^{2q} = \frac{C_q}{N^{D_q(q-1)}}$$

# Previous talk: no multifractal phase on granulated RRG (NL Sigma model)

**But...**

Cuevas, VEK, Ioffe,  
Altshuler (unpublished)



# Gaussian Rosenzweig-Porter RMT: rigorously proven multifractal phase

$$\langle H_{nm} \rangle = 0$$

$N \times N$   
matrix,  
uncorrelated  
random  
entries

Scaling with  
matrix size

$$\sigma = \frac{\lambda^2}{N^\gamma} \ll 1$$

Special diagonal:

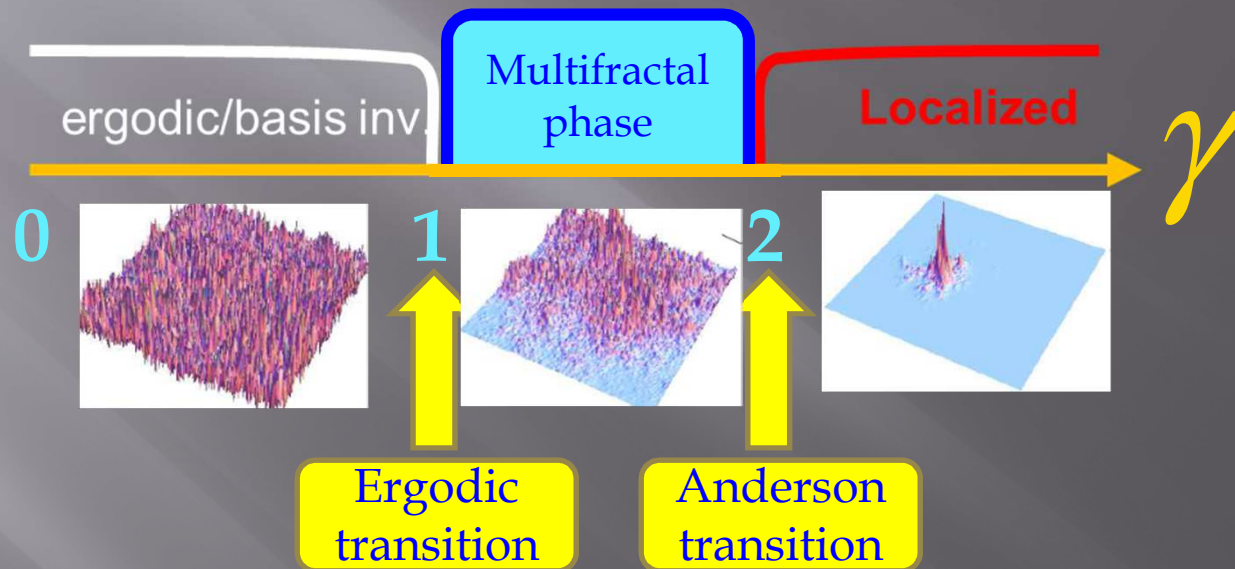
$$\langle H_{nm}^2 \rangle = \begin{pmatrix} 1 & \sigma & \sigma & \sigma \\ \sigma & 1 & \sigma & \sigma \\ \sigma & \sigma & 1 & \sigma \\ \sigma & \sigma & \sigma & 1 \end{pmatrix}$$

Simplest non-invariant RMT



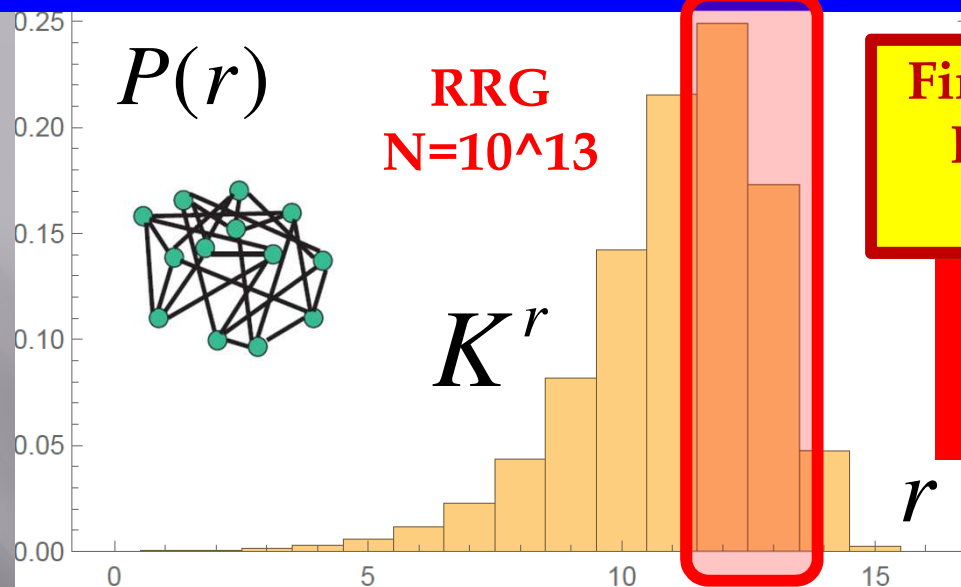
# Multifractal phase in Gaussian RP

V.E.K., I.M. Khaymovich, E. Cuevas, M. Amini,  
New J. Phys., v.17, 12202 (2015)



# Anderson model on RRG and Rosenzweig-Porter ensemble

Two-point Green's function  $G_{0,r} \sim \exp[-\lambda r]$  on a tree is an effective transmission matrix element



Finite fraction of all distances  $r$  on RRG are at a maximal distance  $d = \ln N / \ln K$

Approximation: neglect all other distances!

Only two distribution functions:  $P(\epsilon)$  for diagonal and  $P(G_{0,d})$  for off-diagonal matrix elements  
=Rosenzweig-Porter RMT

$$H_{nn} = \epsilon_n$$

$$H_{nm} = G_{0,d} \equiv G$$

# Log-normal distribution of off-diagonal elements

$$P(G) = \frac{A}{G} \exp \left[ -\frac{\ln^2 (G / G_{typ})}{2 \ln G_{typ}^{-1}} \right]$$

$$G_{typ} \sim N^{-\gamma/2}$$

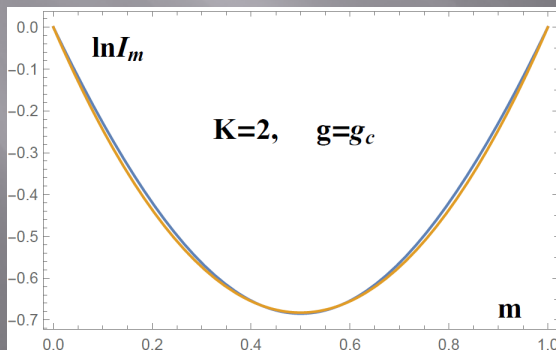
$$\gamma = \frac{2\lambda(W)}{\ln K}$$

Exact for NLσM,  $g \gg 1$

**p-extension:**

$$\tilde{P}(\ln G) \sim \exp \left[ -\ln N \frac{\left( \frac{\ln G}{\ln N} + \frac{\gamma}{2} \right)^2}{p \gamma} \right]$$

$\lambda(W)$  =  
Lyapunov  
exponent on a  
tree with  
branching  $K$



Good for  $K=2$  up to the  
AT point

**RRG -->  
"multifractal" log-normal  
distribution with  $p=1$**

# Role of tails of the lognormal distribution

$$G_{typ} = \exp[\langle \ln G \rangle] \sim N^{\gamma/2}$$

$$\langle G \rangle \sim N^{\gamma_{av}/2}$$

$$\gamma_{av} = (1 - p/2)\gamma$$

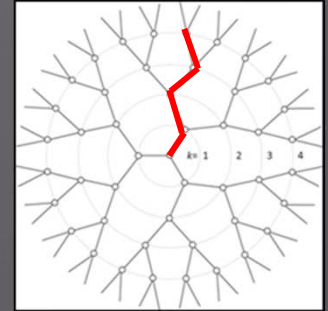
**Averaged off-diagonal matrix elements have different scaling in  $N$  compared to the typical ones: effect of tails in the distribution**

# Imbedded symmetry

For a tree (or 1D system):

$$G_{0,r} = \prod_{n \in \text{path}} G_{nn}$$

$$y = \prod_{n \in \text{path}} G_{nn}^{-1} = G_{0r}^{-1}$$



Basic symmetry  
on a tree

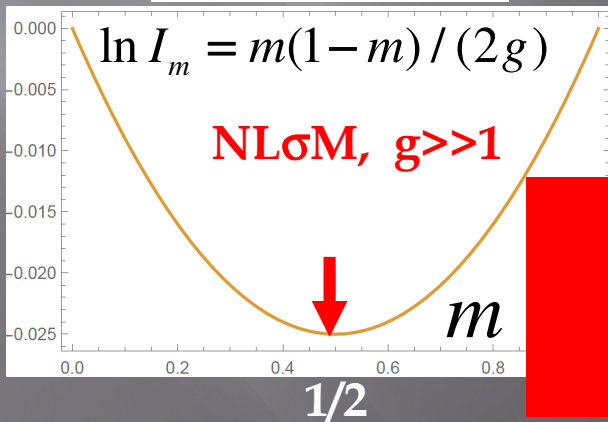
$$F(y) = F(1/y)$$

$$P(1/G) = G^4 P(G)$$

$$(I_m)^r = \int \frac{dy}{y^{2m}} F(y)$$

$$I_m = I_{1-m}$$

$$F(y) = \frac{2}{y} \int_B \frac{dm}{2\pi i} y^{2m} (I_m)^r$$



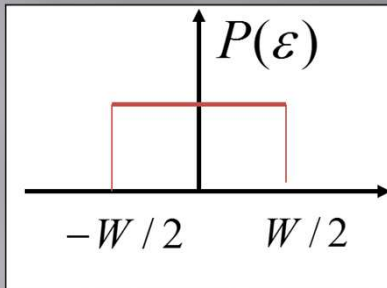
$p=1$   
is a symmetry parameter

$$\tilde{P}(\ln G) \sim \exp \left[ -\frac{\ln N \left( \frac{\ln G}{\ln N} + \frac{\gamma}{2} \right)^2}{p \gamma} \right]$$

# Criterion of localization transition

Localization transition: few sites resonant with a given one (Anderson's criterion):

$$P_{res} = W^{-1} \int_0^W d\omega \int_{\omega=|\varepsilon_n - \varepsilon_m|}^{\infty} dG P(G) \sim 1/N$$



$$N \langle G \rangle_W \sim N^0$$

Truncation of lognormal at  $G \sim W \sim O(1)$

$$\gamma^{(AT)} = \begin{cases} \frac{4}{2-p}, & \text{if } p < 1 \\ 4p, & \text{otherwise} \end{cases}$$

# Criterion of ergodic transition

Breit-Wigner width  $\Gamma$  is of the order of disorder strength  $W$  (Mott's criterion)

$$\Gamma \equiv N \delta_{typ} \sim \sum_j G_{ij}^2 \rho_j \sim \langle G^2 \rangle \delta_{typ}^{-1} \sim W$$

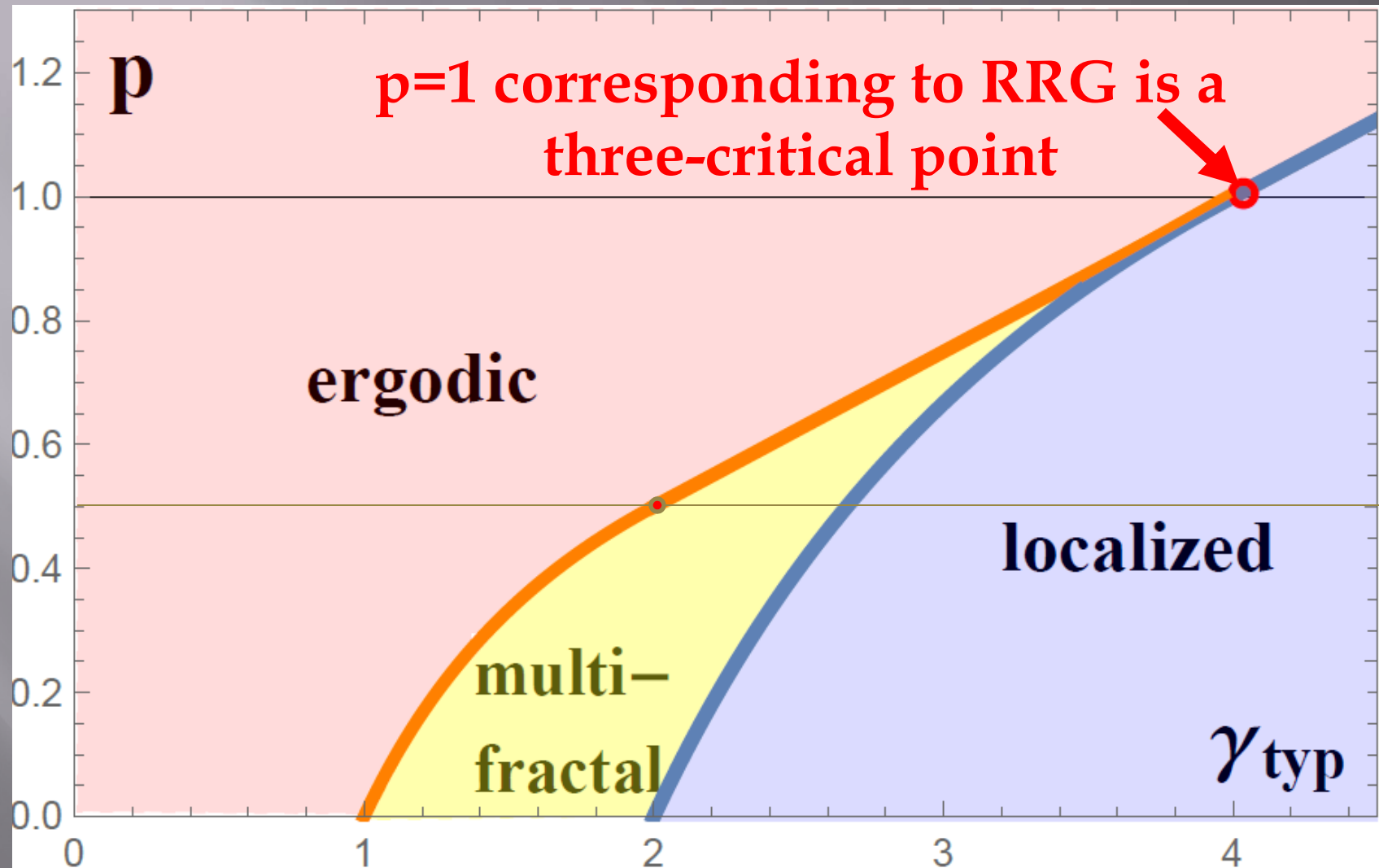
$$N \langle G^2 \rangle_W \sim W^2 \sim N^0$$

Truncation of lognormal at  $G \sim W \sim O(1)$

$G > W$  lead to states at the Lifshits tail which do not contribute to  $\delta_{typ}$

$$\gamma^{(ET)} = \begin{cases} \frac{1}{1-p}, & \text{if } p < \frac{1}{2}, \\ 4p, & \text{if } p > \frac{1}{2} \end{cases}$$

# Phase diagram





# KL (Kullback-Leibler) statistics

$$KL1 = \left\langle \sum_r |\psi_n(r)|^2 \ln \left( \frac{|\psi_n(r)|^2}{|\psi_{n+1}(r)|^2} \right) \right\rangle$$

Sensitive to  
localization transition

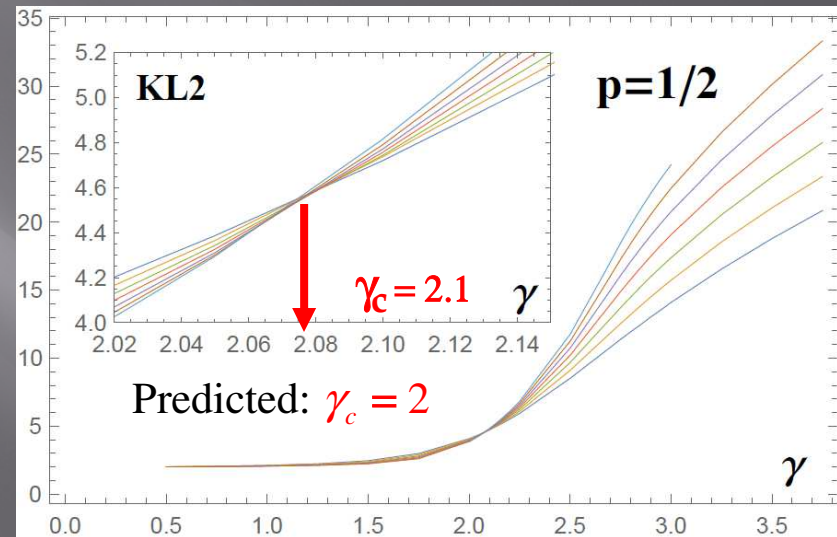
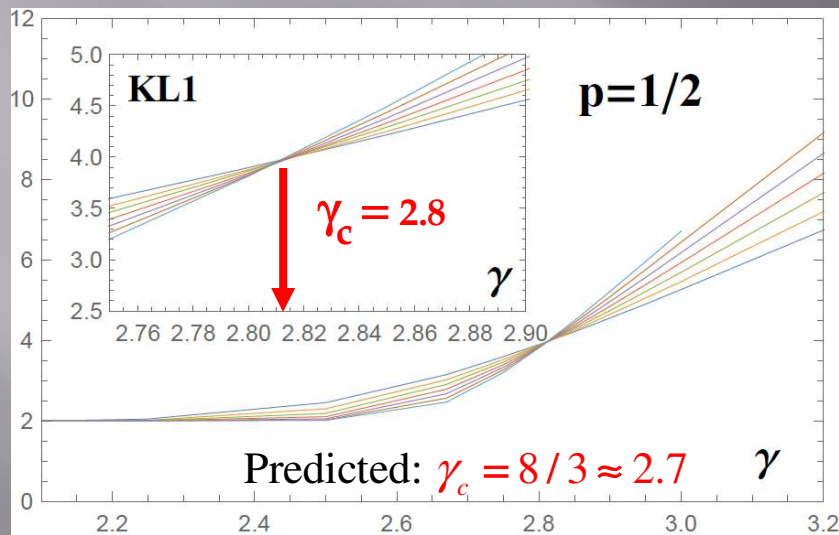
**One and the same disorder realization. Neighboring states strongly correlated in extended phase and uncorrelated in localized phase**

$$KL2 = \left\langle \sum_r |\psi(r)|^2 \ln \left( \frac{|\psi(r)|^2}{|\varphi(r)|^2} \right) \right\rangle$$

Sensitive to ergodic  
transition

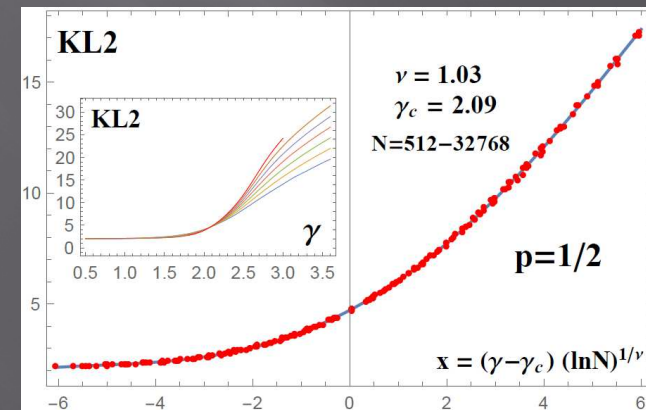
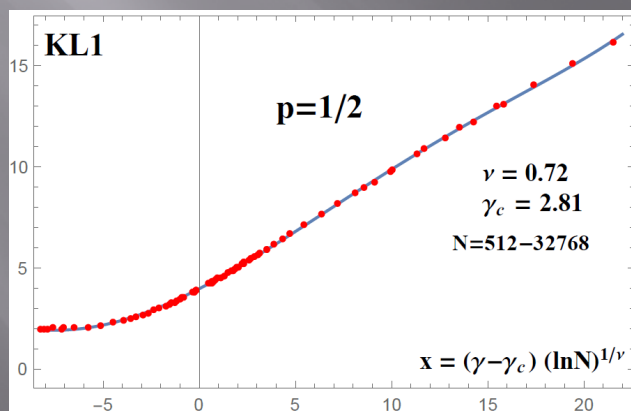
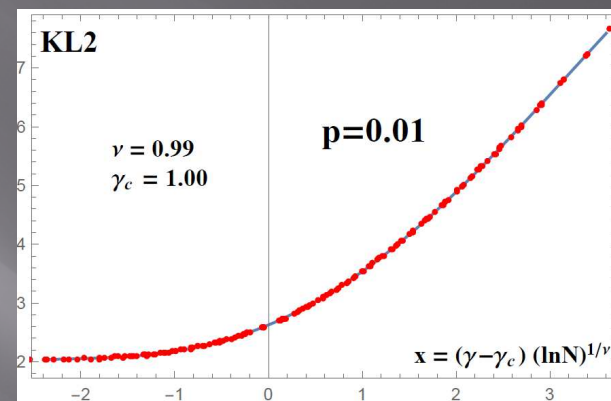
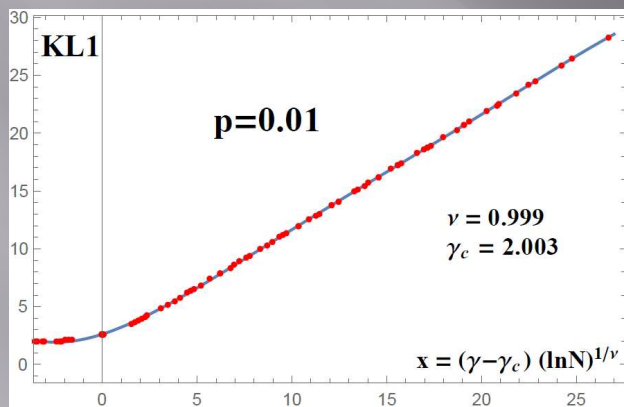
**Different disorder realizations for  $\varphi$  and  $\psi$ . Correlated only in the ergodic phase**

# KL statistics for $p=1/2$

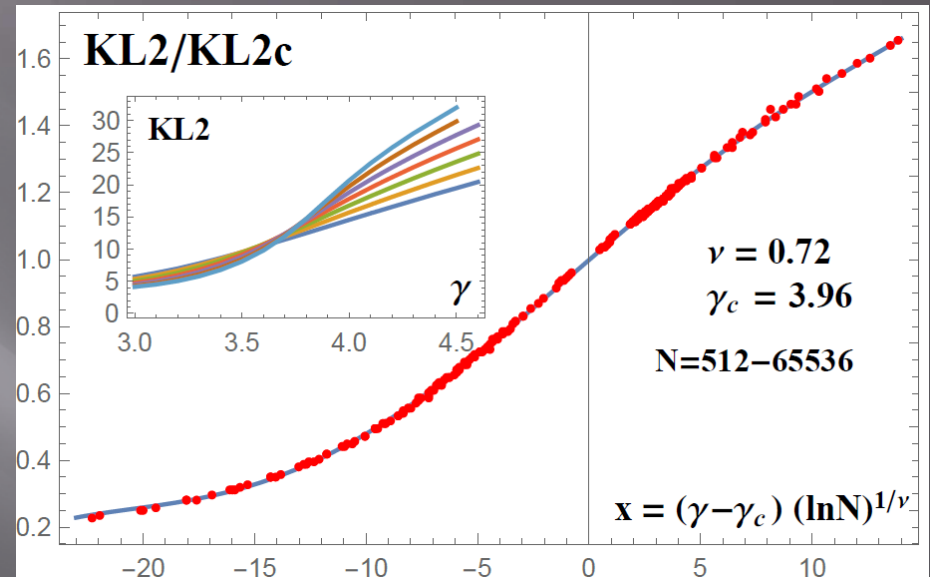
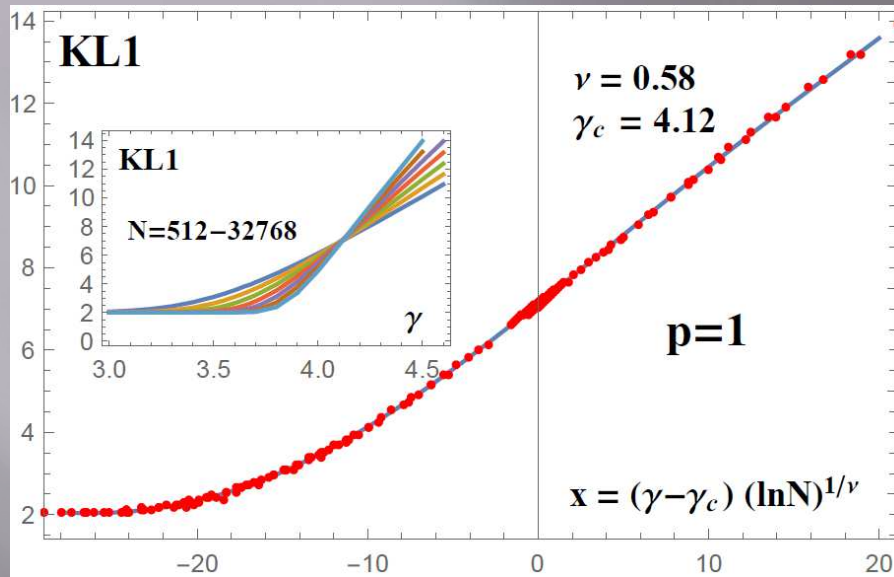


# Data collapse and exponents $\nu$

$$\xi \sim |\gamma - \gamma_c|^{-\nu}$$



# Data collapse and exponents $\nu$ at the three-critical point $p=1$

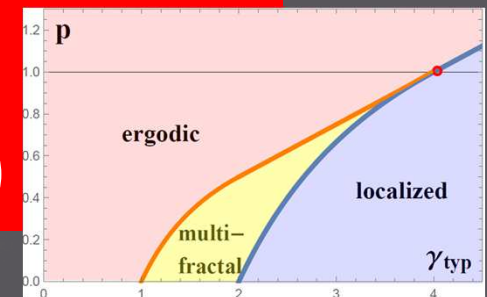


**Merging of two transition points at  $\gamma=4$**

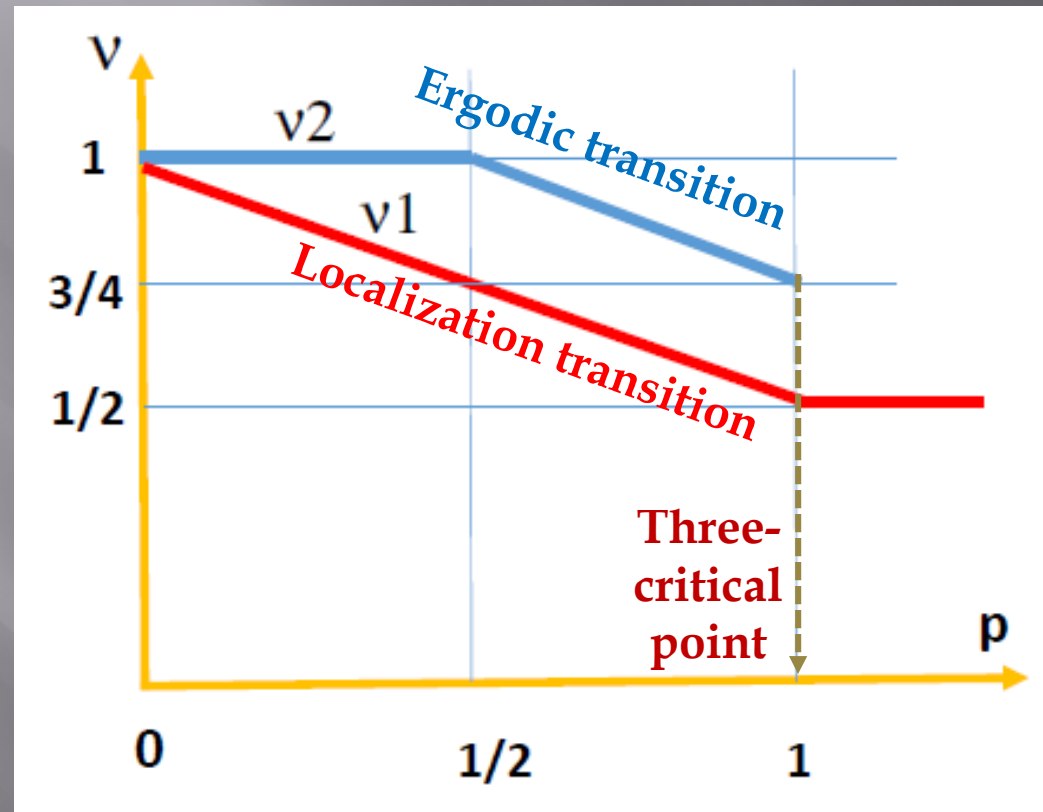
**Two different critical exponents:**

**$\nu=1/2$  and  $\nu=3/4$   
(two critical lengths)**

For a BL case see  
Efetov's  
"Supersymmetry"  
book, 1990



# Conjecture about critical exponents $\nu$

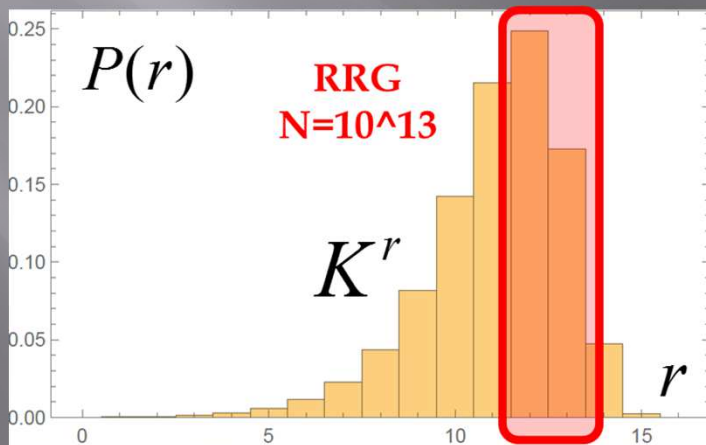


# Self consistency eq. and RP

Mirlin and Fyodorov 1992  
Tikhonov and Mirlin, 2018

$$g_0(Q) = \int e^{-\text{Str} \left[ -2g(Q-Q')^2 + \frac{\pi\eta}{\delta_0} \Lambda Q' \right]} g_0^K(Q') DQ'$$

Zero-dimensional Efetov's supermatrices:  
non-zero space modes are neglected



$$H_{nn} = \varepsilon_n$$

$$H_{nm} = G_{0,d} \equiv G$$

i.i.d. off-diagonal matrix elements

Zero-dimensional  
RP random  
matrices

Equivalent  
APPROXIMATIONS?

# Conclusion

- ▣ Rosenzweig-Porter RMT corresponding to RRG and its  $p$ -extension
- ▣ “Multifractal”, tailed distribution of off-diagonal matrix elements
- ▣ RRG symmetry requirement  $p=1$
- ▣ Criteria of localization and ergodicity
- ▣ Phase diagram for log-normal RP RMT;  $p=1$  is a three-critical point
- ▣ Kullback-Leibler statistics of eigenvectors
- ▣ Numerical characterization of Anderson and Ergodic transitions in logarithmically-normal RP RMT, critical length and exponents  $\nu$